#### **RADIATIVE COOLING**

# Temperature-adaptive radiative coating for all-season household thermal regulation

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The sky is a natural heat sink that has been extensively used for passive radiative cooling of households. A lot of focus has been on maximizing the radiative cooling power of roof coating in the hot daytime using static, cooling-optimized material properties. However, the resultant overcooling in cold night or winter times exacerbates the heating cost, especially in climates where heating dominates energy consumption. We approached thermal regulation from an all-season perspective by developing a mechanically flexible coating that adapts its thermal emittance to different ambient temperatures. The fabricated temperature-adaptive radiative coating (TARC) optimally absorbs the solar energy and automatically switches thermal emittance from 0.20 for ambient temperatures lower than 15°C to 0.90 for temperatures above 30°C, driven by a photonically amplified metal-insulator transition. Simulations show that this system outperforms existing roof coatings for energy saving in most climates, especially those with substantial seasonal variations.

n countries such as the United States, ~39% of the total energy consumption is in buildings (1). For the residential housing energy portion, ~51% is consumed for heating and cooling to maintain a desirable indoor temperature (~22°C) (2). In contrast to most temperature regulation systems, which require external power input, the mid-infrared (IR) atmospheric transparency window ("sky window") allows thermal radiation exchange between terrestrial surfaces and the 3 K outer space, thus opening a passive avenue for thermal radiative cooling of buildings. This method to cool an outdoor surface such as a roof has been extensively studied in the past (3-6). It is now advanced by the development of daytime radiative cooling (7-13) using materials with low solar absorptance and high thermal emittance in the form of thin films (8), organic paints (10, 14), or structural materials (11).

Past research on daytime radiative cooling, while successful in reducing cooling energy consumption, typically used materials with fixed, cooling-optimized properties, which efficiently emit thermal radiation even when the temperature of the surface is lower than de-

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We took a different, holistic approach by designing and fabricating a mechanically flexible coating structure to minimize total energy consumption through the entire year. This temperature-adaptive radiative coating (TARC) automatically switches its sky-window emittance to 0.90 from 0.20 when the surface temperature rises above ~22°C, a practical threshold not previously available. Our TARC delivers high radiative cooling power exclusively for the high-temperature condition (Fig. 1A). We also optimized the solar absorptance at ~0.25 (solar reflectance = 0.75) for all-season energy saving in major US cities (fig. S7). Our TARC demonstrates effective surface temperature modulation in an outdoor test environment. We performed extensive simulations based on the device properties and the climate database, which show advantages of TARC over existing roof coating materials in energy savings for most US cities in different climate zones (Fig. 1C). The energy savings by TARC not only bring economic benefits but also contribute to environmental preservation by reducing greenhouse gas emissions.

We developed the TARC based on the wellknown metal-insulator transition (MIT) of the strongly correlated electron materials  $W_x V_{1-x} O_2$ (30-32), and the transition temperature  $(T_{\rm MIT})$ is tailored to ~22°C by setting the composition x at 1.5% (33). We embedded a lithographically patterned two-dimensional array of thin  $W_x V_{1-x} O_2$  blocks in a BaF<sub>2</sub> dielectric layer that sits on top of an Ag film (Fig. 2A). In the insulating (I) state of  $W_x V_{1-x} O_2$  at  $T < T_{MIT}$ , the material is largely transparent to the infrared (IR) radiation in the 8- to 13-um sky spectral window, so this sky-window IR radiation is reflected by the Ag mirror with little absorption (34). By contrast, the  $W_x V_{1-x} O_2$  becomes highly absorptive in the sky window when it switches to the metallic (M) state at  $T > T_{MIT}$ (34). The absorption is further amplified by the designed photonic resonance with adjacent  $W_x V_{1-x} O_2$  blocks as well as with the bottom Ag layer through the ¼-wavelength cavity. The ¼-wavelength cavity structure induces Fabry-Perot resonance and was used in previous work to enhance thermal emission (21, 23). According to Kirchhoff's law of radiation (35), the sky-window emittance equals the skywindow absorptance and switches from low to high when the temperature exceeds  $T_{\rm MIT}$ . Consequently, strong sky-window radiative cooling is turned on in operation exclusively at high temperatures, leaving the system in the solar-heating or keep-warm mode at low temperatures. Details on the fabrication process and structural parameters are found in the supplementary materials (36) (fig. S1).

Our fabricated TARC has high flexibility for versatile surface adaption, as well as a microscale structure consistent with the design (Fig. 2B). We examined the emittance switching over the entire sample using a thermal infrared (TIR) camera (Fig. 2C). We imaged the TARC surface together with two reference samples having similar thicknesses but constant low thermal emittance (0.10, copper plate) or constant high thermal emittance (0.95, black tape), respectively. Although the thermal emission of the reference samples appeared to not be strongly temperature sensitive from 20 to 30°C, the TARC showed a marked change, corresponding to the switch in sky-window emittance at the MIT around 22°C.

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We measured the spectral properties of the TARC by a UV-visible-NIR spectrometer and Fourier transform infrared spectroscopy (FTIR) for the solar and TIR wavelength regimes, respectively (Fig. 2D). The solar absorptance (A, 0.3 to 2.5 µm) is ~0.25, and the sky-window emittance ( $\varepsilon_w$ , 8 to 13 µm) is ~0.20 in the I state and ~0.90 in the M state, consistent with theoretical simulations and other characterization results (fig. S2 and fig. S3).

The emittance switching of the TARC enables deep modulation of radiative cooling power in response to ambient temperature, which we first measured in vacuum (Fig. 3A). We suspended a heater membrane by thin strings in a vacuum chamber, which was cooled with dry ice to ~-78°C to minimize radiation from the chamber walls. We attached a piece of Al foil with  $\epsilon_{Al}\approx$  0.03 or a TARC of the same size to the top of the heater in two separate measurements. At each stabilized sample temperature T, the heating powers needed for the two coating scenarios are denoted as  $P_{Al}$ (T) and  $P_{\text{TARC}}(T)$ , respectively. The cooling flux (power per area  $\mathcal{A}$ ) contributed by the TARC was calculated as  $P'_{\text{cool}}(T) = [P_{\text{TARC}}(T) P_{\rm Al}(T)]/\mathcal{A}$ . We used the Al foil reference to calibrate background heat loss from thermal conduction through the strings. We plotted the calibrated cooling power (Fig. 3B), which shows an abrupt increase in  $P''_{cool}(T)$  when T rises above the MIT temperature.  $P''_{cool}(T)$ measurements in the I state and M state are well fitted by the Stefan-Boltzmann radiation law, with values of sky-window  $\varepsilon_w$  extracted to be  $\sim 0.20$  and  $\sim 0.90$ , respectively. consistent with the spectrally characterized results (Fig. 2D). We considered and corrected the effect of radiation from the chamber wall (~-78°C) for the calibration. We introduced a constant factor of  $\gamma ~(\approx 0.7)$  to account for the difference between the vacuum and ambient measurement conditions (details in fig. S4) (36).

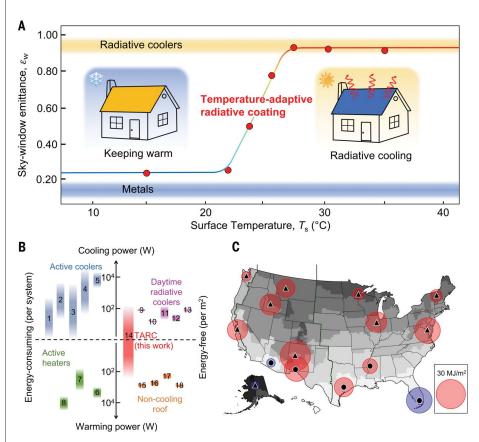
We demonstrated the actual outdoor performance of the TARC (Fig. 4). We recorded the surface temperatures ( $T_s$ ) of the TARC, together with a dark roof coating product (Behr no. N520, asphalt gray) and a cool (white) roof coating product (GAF RoofShield white acrylic), over 24 hours on a sunny summer day on a rooftop in Berkeley, California, with a careful design of the measurement system to minimize the effects of artifacts (fig. S5).

From 00:00 to 09:00 local daylight time (LDT), when the ambient temperature was below  $T_{\rm MIT}$ , the TARC was 2°C warmer than the two reference roof coatings, arising from the low sky-window emittance ( $\varepsilon_{\rm w} = 0.20$ ) of the TARC in the I state and thus a lower radiative cooling power than the references ( $\varepsilon_{\rm w} = 0.90$ ). The 2°C temperature elevation is consistent with adiabatic simulation results based on these nominal emittance values and the local

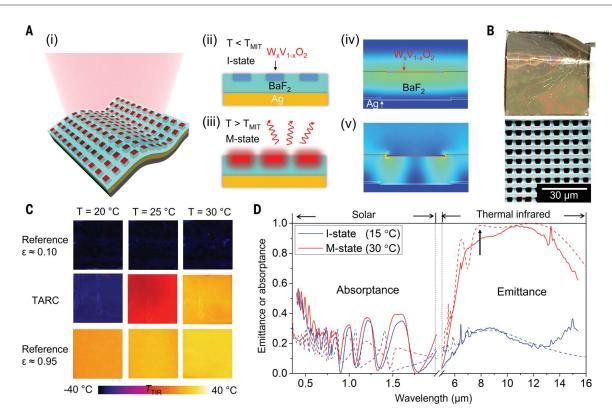
weather database [see the supplementary materials (36), note A, section I]. From 09:00 to 13:00 LDT, when the samples were in direct sunlight,  $T_s$  was dominated by the solar absorption in balance with radiative cooling and air convection, and the differences between the samples agree with the simulated results assuming the solar absorptance A to be 0.15, 0.25, and 0.70 for the white roof coating, TARC, and the dark roof coating, respectively. After 13:00 LDT, we erected a shield to intentionally block direct solar radiation to the surface of the samples. This imitates the scenario of a cloud blocking the sun but with the rest of the sky mostly clear. We quickly observed a convergence of the  $T_{\rm s}$  curves for all three samples, an indication that the thermal emittance of the TARC in the M state is close to that of the two references (0.90). This condition persisted for a few hours until  $T_s$  started to drop below  $T_{\rm MIT}$  = 22°C. After this point, TARC grew warmer than the two references, with a

final temperature difference of ~2°C, similar to the 00:00 to 09:00 LDT period. This indicates that the TARC switched to the low-emittance I state. The 24-hour outdoor experiments demonstrate the emittance switching and resultant temperature regulation by TARC. Although the white roof coating shows an advantage over TARC in thermal management in summer daytime and under solar radiation (Fig. 4A), the TARC regulates the roof temperature closer to the heating and cooling setpoints (22 and 24°C) than the white roof coating for almost all of the other conditions, including daytime in other seasons and all of the nighttime (fig. S6). From an all-year-round perspective, the TARC demonstrates superiority compared with regular roof coatings in terms of source energy saving.

To directly compare their ambient condition cooling fluxes  $(P'_{cool,amb})$ , we heated the TARC and the white roof coating to the air temperature with the direct solar radiation



**Fig. 1. TARC and its benefits for household thermal regulation.** (**A**) Basic property of TARC in sky-window (8 to 13  $\mu$ m) emittance modulation and schematics for temperature management when used as a household roof coating. The data points are the measured sky-window emittances of a TARC. The two color bands represent the temperature-independent thermal emittance of metals and radiative coolers. (**B**) TARC compared with other thermal regulation systems, highlighting the unique benefit of TARC of being simultaneously energy-free and temperature adaptive (details in table S1). (**C**) SCSES<sub>min</sub> of TARC compared with other existing roof-coating materials for different cities representing the 15 climate zones in the United States. Red and blue circles indicate positive and negative SCSES<sub>min</sub> values, respectively. The values are scaled to the area of the circles. Representation of the triangle and circle icons is explained in the materials and methods (subsection, "Projection of energy savings") (36).



# Fig. 2. Basic properties of TARC with experimental characterization. (A) Schematics of the structure (i), materials composition and working mechanism (ii and iii) of the TARC. Subpanels (iv) and (v) show the simulated distribution of electric field intensity below and above the transition temperature, respectively, when electromagnetic waves with a wavelength of 7.8 $\mu$ m were normally incident on the TARC structure. (B) Photograph (2 cm × 2 cm) and false-color scanning electron microscope image of TARC showing high flexibility and structural consistency with the design. (C) TIR images of TARC compared with those of two conventional

materials (references) with constantly low or high thermal emittance showing the temperature-adaptive switching in thermal emittance of TARC. (**D**) Solar spectral absorptance and part of the thermal spectral emittance of TARC at a low temperature and a high temperature, measured by a UV-visible-NIR spectrometer with an integrating sphere and an FTIR spectrometer, respectively. Measurements (solid curves) show consistency with theoretical predictions (dashed curves). The arrow at 7.8  $\mu$ m denotes the wavelength where the distribution of electric field intensity shown in subpanels (iv) and (v) of (A) are simulated.

blocked.  $P'_{\text{cool.amb}}$  refers to the net cooling flux from the surface—namely, the thermal radiative heat loss flux minus the absorbed diffuse solar irradiance. We plotted the  $P'_{\text{cool.amb}}$ values that we obtained at a low and a high air temperature (Fig. 4B). The TARC exhibits a clear switching of  $P'_{\text{cool.amb}}$  by a factor over five across the MIT. This behavior is in stark contrast to the nearly constant  $P'_{\text{cool.amb}}$  around 120 W/m<sup>2</sup> for the shaded white roof coating, which is consistent with values (90 to 130 W/m<sup>2</sup>) reported in literature for roofs surfaced with daytime radiative cooling materials (*5, 9, 10*).

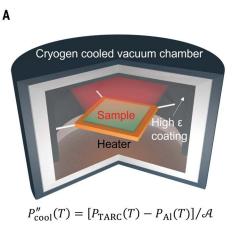
We performed extensive numerical simulations to analyze the performance of TARC in household energy saving for the US cities from an all-season perspective (*36*). We show the simulated results (Fig. 4C) for Berkeley where the measurements (Fig. 4, A and B) were performed. We calculated an hour-month map of  $T_{\rm s}$  using a local weather file (*37*), laying the basis for estimation of energy saving. We assumed heating and cooling setpoints  $T_{\rm set,heat} =$ 22°C and  $T_{\rm set,cool} = 24$ °C (*38*), and approx-

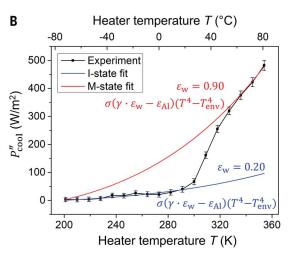
imated that the building will need heating when  $T_s < T_{set,heat}$  and require cooling when  $T_{\rm s} > T_{\rm set, cool}$ . We used past simulations of coolroof energy savings to predict potential spaceconditioning source energy savings (SCSES) per unit roof area attainable by using TARC in place of roofing materials that have static values of solar absorptance and thermal emittance (36). The figure of merit of TARC is represented by SCSES<sub>min</sub>, the minimum value of SCSES found over all existing conventional roofing materials, which have constant values of  $A_{\rm ref}$  and  $\varepsilon_{\rm ref}$  (Fig. 4C, dashed boxes). We mapped SCSES<sub>min</sub> for cities representing the 15 US climate zones (Fig. 1C). This figureof-merit map shows that TARC provides clear, positive annual space-conditioning source energy savings relative to existing roof coating materials in most major cities, except for climates that are constantly cold (such as Fairbanks) or hot (such as Miami) throughout the year. It highlights the advantage of TARC, especially in climate zones with wide temperature variations, day to night or summer to winter. For example, we estimate that for a single-family home in Baltimore, Maryland, built before 1980, modeled with roof assembly thermal insulance 4.3 m<sup>2</sup>/(K·W), gas furnace annual fuel utilization efficiency 80%, and air conditioner coefficient of performance 2.64 (*38*), SCSES<sub>min</sub> is 22.4 MJ/(m<sup>2</sup>·y), saving 2.64 GJ/y based on a roof area of 118 m<sup>2</sup>. We also calculated the source energy saving of TARC as a function of its solar absorptance (fig. S7), showing that the actual solar absorptance of TARC is close to the optimal value for major US cities.

The TARC could be readily upgraded for heavy-duty outdoor applications by coating it with a thin polyethylene (PE) membrane, which is nontoxic, hydrophobic, and transparent both in the visible and thermal IR regions. While protecting the TARC from contacting the dust and moisture in complex environments, the PE coating has little impact on the thermal modulation performance (fig. S9). Polymer imprinting instead of photolithography could also be used to more easily produce the material for large scale application. By embedding VO<sub>2</sub> particles in layered PE

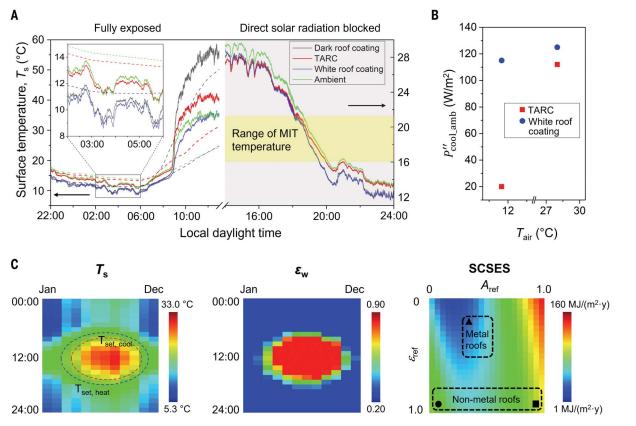
#### Fig. 3. Characterization of intrinsic radiative cooling power of TARC in a cryogenic vacuum chamber. (A) Schematics of the experimental setup showing a thin heater membrane covered by either a TARC or an Al foil and suspended in a cryogen-cooled vacuum chamber. The Al foil reference is used to cancel the effect of thermal loss through conduction. (B) Calibrated

conduction. (**B**) Calibrated experimental cooling flux (power/ area) of TARC as a function of temperature in vacuum (the black data line). Fitting of  $P_{cool}^{"}(T)$  at I and M states by the Stefan-





Boltzmann radiation law gives  $\varepsilon_w$  values of 0.20 and 0.90, respectively.



**Fig. 4. Characterization of TARC in an outdoor environment.** (**A**) Surface temperature of TARC, a commercial dark roof coating (A = 0.70,  $\varepsilon_w = 0.90$ ), and a commercial white roof coating (A = 0.15,  $\varepsilon_w = 0.90$ ) in an open-space outdoor environment recorded over a day-night cycle. The measurement was taken on 5 July 2020, in Berkeley, California (37.91°N, 122.28°W). The solid and dashed curves are experimental data and simulation results based on a local weather database (*37*), respectively. Measurements starting from 14:00 LDT were performed with the direct solar radiation blocked. Temperature observed after sunset show clear signs

of the TARC shutting off thermal radiative cooling as its surface ambient temperature falls below  $T_{MIT}$ . (**B**) Measured ambient cooling power of TARC and white roof coating with direct solar radiation blocked in the outdoor environment. (**C**)  $T_{\rm s}$  and the corresponding  $\varepsilon_{\rm w}$  mapping of TARC over 24 hours and the full year for Berkeley. Also shown are the SCSES of TARC compared with all other materials with fixed solar absorptance ( $A_{\rm ref}$ ) and fixed thermal emittance ( $\varepsilon_{\rm ref}$ ). The icons in the SCSES map correspond to those used in Fig. 1C, denoting the radiative parameters (A,  $\varepsilon_{\rm w}$ ) of the strongest rival to TARC in source energy savings for the local climate (*36*).

membranes, we estimated the multilayered metamaterial to achieve comparable modulation performance ( $\Delta \varepsilon_w > 0.8$ ) as the TARC we presented and would be producible in a roll-to-roll fashion (figs. S10 and S11). Roll-to-

roll manufacturing of PE-based TARC would be beneficial because of its high scalability, low cost (9), and the fact that it is free from the liquid evaporation process in fabrication (39). The PE layer can be also replaced by other organic or inorganic materials with negligible optical loss in the wavelength ranges of both solar irradiation and IR atmospheric transparency window, so that the TARC technology can be designed specifically to be endurable in different environmental conditions.

We developed a mechanically flexible, energyfree TARC for intelligent regulation of household temperature. Our system features a thermally driven metal-insulator transition in cooperation with photonic resonance, and demonstrates self-switching in sky-window thermal emittance from 0.20 to 0.90 at a desired temperature of ~22°C. These attractive properties enable switching of the system from the radiative cooling mode at high temperatures to the solar-heating or keep-warm mode at low temperatures in an outdoor setting. For most cities in the United States, our simulations indicate the TARC may outperform all conventional roof materials in terms of cutting energy consumption for households.

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directors of the Cool Roof Rating Council (CRRC) and a paid consultant to the CRRC. K.T., K.D., J.L., and J.W. are inventors of a provisional patent application related to this work. The authors declare that they have no competing interests. **Data and materials availability:** All data required to evaluate the conclusions in the manuscript are available in the main text or the supplementary materials.

#### SUPPLEMENTARY MATERIALS

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## Science

## Temperature-adaptive radiative coating for all-season household thermal regulation

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#### A passive turnoff

Passive radiative cooling technology uses the infrared atmospheric window to allow outer space to be a cold sink for heat. However, this effect is one that is only helpful for energy savings in the warmer months. Wang *et al.* and Tang *et al.* used the metal-insulator transition in tungsten-doped vanadium dioxide to create window glass and a rooftop coating that circumvents this problem by turning off the radiative cooling at lower temperatures. Because the transition is simply temperature dependent, this effect also happens passively. Model simulations suggest that these materials would lead to energy savings year-round across most of the climate zones in the United States. —BG

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## Supplementary Materials for

## Temperature-adaptive radiative coating for all-season household thermal regulation

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## Nomenclature

## Abbreviations

Ag	silver		
$BaF_2$	barium fluoride		
DI	deionized		
FTIR	Fourier transform infrared		
HF	hydrofluoric acid		
HNO <sub>3</sub>	nitride acid		
HVAC	heating, ventilation, and air conditioning		
IR	infrared		
LDT	local daylight time		
MIT	metal-insulator transition		
$N_2$	nitrogen		
NIR	near-infrared		
O <sub>2</sub>	oxygen		
PE	polyethylene		
PLD	pulsed laser deposition		
SCSES	space conditioning source energy saving		
Si	silicon		
$SF_6$	sulfur hexafluoride		
TARC	temperature-adaptive radiative coating		
TIR	thermal infrared		
TMY3	typical meteorological year 3		
U.S.	United States		
UV	ultraviolet		
$VO_2$	vanadium dioxide		
$V_2O_5$	vanadium pentoxide		
WO <sub>3</sub>	tungsten trioxide		
$W_x V_{1-x} O_2$	tungsten-doped vanadium oxide		
XeF <sub>2</sub>	xenon difluoride		

## English Symbols

Α	solar absorptance
$A_{\rm ref}$	solar absorptance of reference materials
$B(\lambda,T)$	spectral radiance emitted by a black body at absolute temperature $T$
С	annual source energy uses for air conditioning
CF	cloud coverage factor
D <sub>c</sub>	cooling degrees
$D_{\rm h}$	heating degrees
$\Delta D_{\rm c}$	reduction in cooling degrees
$\Delta D_{ m h}$	reduction in heating degrees
$\Delta D_{\rm c,TARC}$	reduction in cooling degrees by TARC
$\Delta D_{ m h,TARC}$	reduction in heating degrees by TARC
F	annual fan source energy use
<i>F</i> <sub>c</sub>	annual cooling fan source energy use
<i>F</i> <sub>h</sub>	annual heating fan source energy use
Н	annual source energy uses for gas heating
Ι	solar irradiance
$I_{\rm s}(\lambda)$	solar spectral irradiance
P <sub>Al</sub>	measured heater power for Al surface in vacuum
P <sub>TARC</sub>	measured heater power for TARC surface in vacuum
$P_{ m cool}''$	cooling flux (power/area) of TARC in vacuum
$P_{\rm cool\_amb}''$	cooling flux (power/area) of TARC in ambient
q <sub>c</sub>	convective heat loss
$q_{ m LW}$	long-wave thermal radiative loss
$q_{\rm SW}$	short-wave solar absorption
R	solar reflectance
$r(\lambda)$	spectral reflectance
$T_{ m MIT}$	metal-insulator-transition temperature
T <sub>a</sub>	air temperature

T <sub>d</sub>	dew point temperature
T <sub>s</sub>	surface temperature
$T_{ m set,\ cool}$	setpoint temperature for cooling
T <sub>set, heat</sub>	setpoint temperature for heating
$\Delta S$	annual space-conditioning source energy savings
$\Delta S_{\rm c}$	annual space cooling source energy savings
$\Delta S_{\rm h}$	annual space heating source energy savings
S <sub>TARC</sub>	space-conditioning source energy consumption of TARC
v	wind speed

## Greek Symbols

$\alpha_{\rm c}$	linear fitting coefficient relating $\Delta S_c$ to $\Delta D_c$
$lpha_{ m h}$	linear fitting coefficient relating $\Delta S_h$ to $\Delta D_h$
$a(\lambda)$	spectral absorptance
${\mathcal A}$	surface area
$\beta_{\rm c}$	linear-fit coefficient relating $F_c$ to C
$eta_{ m h}$	linear-fit coefficient relating $F_{\rm h}$ to $H$
γ	calibration factor for thermal emittance between vacuum and ambient conditions
$\varepsilon_{ m ref}$	sky-window thermal emittance of reference materials
$\mathcal{E}_{\mathrm{W}}$	sky-window thermal emittance
$\epsilon_{\rm vac}$	thermal emittance in vacuum
ε <sub>clear_sky</sub>	sky emissivity in a clear-sky model
$\varepsilon_{ m sky}$	sky emissivity with cloud correction
$\varepsilon(\lambda)$	spectral emittance

## **Materials and Methods**

## **Preparation of the TARC**

675 µm-thick Si wafers were first covered with an approximately 2 µm thick polyimide film (PI-2545, HD MicroSystems LLC) via spin coating, which was then cured in a N<sub>2</sub>-filled oven at 425 °C for 60 minutes. The polyimide film acts as an etching protection layer for the final transfer process.  $W_xV_{1-x}O_2$  thin films were grown on the polyimide layer using pulse laser deposition (PLD). The PLD target was prepared by mixing WO<sub>3</sub> and V<sub>2</sub>O<sub>5</sub> powders with a W:V atomic ratio at 1.5%, then made into 2.5 cm diameter round discs with a hydraulic press. All  $W_xV_{1-x}O_2$  thin films were deposited in a 5 mTorr O<sub>2</sub> environment at 500 °C substrate temperature, and the PLD laser energy was set at 321 mJ with 10 Hz pulse frequency. 70 nm of  $W_xV_{1-x}O_2$  was grown at a rate of 6 nm/min, followed by a post-deposition anneal at 500 °C for 30 mins in the same 5 mTorr O<sub>2</sub> environment. The metamaterials patterns were made with standard photolithography, combined with etching of  $W_xV_{1-x}O_2$  by SF<sub>6</sub> + O<sub>2</sub> in a plasma etching system. After removing the photoresist with acetone and O<sub>2</sub> plasma, 1.5 µm thick BaF<sub>2</sub> and 100 nm thick Ag layers were grown sequentially on top via thermal evaporation. The growth rates of BaF<sub>2</sub> and Ag were controlled at 20 Å/s and 2 Å/s, respectively.

In the transfer process, a piece of 0.06 mm thick single-sided sticky cellophane packaging tape was first carefully applied to fully cover the surface, where the Ag layer was stuck to the adhesive side without any residual air bubbles. An initial Si substrate removal process was performed in a HF + HNO<sub>3</sub> solution, mixed by aqueous HF (49% weight percentage) and HNO<sub>3</sub> solution (68% weight percentage) with a volume ratio of 10:1. The samples were taken out and rinsed with DI water to stop the initial etching when the etchant starts to touch down on the polyimide layer. A XeF<sub>2</sub> dry etching process was then carried out to clean off the residue Si. In the final step, the polyimide protection layer was removed by O<sub>2</sub> plasma at 100 mTorr O<sub>2</sub> pressure and 200 W plasma power for about 11 mins.

## Spectrally resolved measurements

Thermal spectral reflectance at normal incidence,  $r(\lambda, T)$ , was characterized by a Nicolet iS50 FTIR spectrometer and Nicolet Continuum microscope over the spectrum 5-15 µm. The objective lens was 32× with 0.65 numerical aperture. A blade aperture of 100 µm × 100 µm was used to select the area of interest. All reflection spectra were normalized to the reflection spectrum of a 300 nm thick gold film. The temperature of the samples (15-50 °C) was controlled by a customized closed-loop thermal stage, connected to a Lakeshore 321 temperature controller. Kirchhoff's law of radiation states that in thermodynamic equilibrium, spectral emittance  $\varepsilon(\lambda, T)$  equals spectral absorptance  $a(\lambda, T)$ . Since the TARC was essentially opaque from 5 to 15 µm, its thermal spectral emittance in this range was computed as  $\varepsilon(\lambda, T) = a(\lambda, T) = 1 - r(\lambda, T)$ . Near normal-hemispherical solar spectral reflectance,  $r(\lambda)$ , was measured from 300 to 2,500 nm with an Agilent Cary 5000 UV-vis-NIR spectrometer equipped with an Internal Diffuse Reflectance Accessory (DRA-2500), which collects both specular and diffuse reflections. Solar spectral absorptance was computed as  $a(\lambda) = 1 - r(\lambda)$  since the film was essentially opaque to sunlight.

The solar absorptance A and sky-window thermal emittance  $\varepsilon_s$  can be calculated from the corresponding spectral data by:

$$A = \left(\int_{0.3 \,\mu\mathrm{m}}^{2.5 \,\mu\mathrm{m}} I_{\mathrm{s}}(\lambda) \,a(\lambda) \,d\lambda\right) / \left(\int_{0.3 \,\mu\mathrm{m}}^{2.5 \,\mu\mathrm{m}} I_{\mathrm{s}}(\lambda) \,d\lambda\right)$$
$$\varepsilon_{\mathrm{w}}(T) = \left(\int_{8 \,\mu\mathrm{m}}^{13 \,\mu\mathrm{m}} B(\lambda) \,\varepsilon(\lambda,T) \,d\lambda\right) / \left(\int_{8 \,\mu\mathrm{m}}^{13 \,\mu\mathrm{m}} B(\lambda) \,d\lambda\right)$$

where  $I_s(\lambda)$  is the solar spectral irradiance, and  $B(\lambda)$  is the spectral radiance of a black body emission.

### Thermal infrared imaging and analysis

TIR images were captured by a FLIR ONE infrared (IR) camera working at a wavelength range of 8-13  $\mu$ m (same as the sky window). To minimize the reflection from the camera and the surroundings, the default viewing angle was set as 15° instead of normal incident direction, and the experiments were performed in an open-area, outdoor environment under a clear (cloud-free) sky. When taking TIR images, the camera measures the incident TIR radiation, and then gives the temperature reading ( $T_{\text{TIR}}$ ) assuming a constant thermal emittance for the target (e.g., 0.90, the default setting of the camera).

## Simulation of device properties

The spectral absorptance of TARC is numerically calculated using COMSOL Multiphysics, with all the geometric parameters matching the original design. Material properties in IR and visible ranges are from Refs. (40-43) and (44-46), respectively. Note that to better predict the actual TARC performance in the 5-15  $\mu$ m range, the imaginary part of BaF<sub>2</sub> permittivity is slightly increased by  $\Delta \varepsilon_i = 0.022\lambda - 0.049$  (unit of  $\lambda$  is  $\mu$ m) by fitting.

### Characterization of thermal emittance modulation in vacuum condition

In the measurement setup shown in Fig. 3A, the power of the heater at the equilibrium surface temperature *T* for the Al foil surface is denoted as  $P_{Al}(T)$ , and the power for the TARC surface is denoted as  $P_{TARC}(T)$ .  $P_{Al}(T)$  and  $P_{TARC}(T)$  are:

$$P_{AI}(T) = \sigma \mathcal{A} \varepsilon_{AI} (T^4 - T_{env}^4) + C(T)$$
$$P_{TARC}(T) = \sigma \mathcal{A} \varepsilon_{TARC,v} (T^4 - T_{env}^4) + C(T)$$

Here  $\sigma$  is the Stefan-Boltzmann constant,  $\mathcal{A}$  is the area of the sample surface,  $\varepsilon_{Al}$  and  $\varepsilon_{TARC,v}$  are effective thermal emittance of Al and TARC in vacuum conditions,  $T_{env}$  is the temperature of the environment (chamber wall), C(T) represents ancillary heat loss power from other radiative source and thermal conduction, and all temperatures are absolute. The ancillary loss is temperature dependent and unknown, but is reasonably assumed to be the same for the surface conditions of the two materials. The inner walls of the vacuum chamber can be treated as a black body due to two reasons: (1) High-emittance tape was used to cover the inner walls (see Fig. S12); (2) The sample area (16 cm<sup>2</sup>) is much smaller than the total area of the inner walls (1,500 cm<sup>2</sup>) (47). Therefore, the cooling flux (power/area) contributed by the TARC, denoted as  $P''_{cool}(T)$ , can be calculated as:

$$P_{\text{cool}}''(T) = [P_{\text{TARC}}(T) - P_{\text{Al}}(T)] / \mathcal{A} = \sigma \left(\varepsilon_{\text{TARC},v} - \varepsilon_{\text{Al}}\right) \left(T^4 - T_{\text{env}}^4\right)$$

In this equation,  $\sigma$  is known,  $\varepsilon_{Al}$  is approximated as 0.03, and  $T_{env}$  is equal to the temperature of the cryogen (195 K, or -78 °C). Therefore,  $\varepsilon_{TARC,v}$  at the I state and the M state can be obtained by fitting  $P''_{cool}(T)$  with T at two branches before and after MIT. The error bars in Fig. 3B come from the uncertainty of  $\varepsilon_{Al}$  (taken as 0.02) and a systematic error  $\epsilon_s$  (taken as 8.0 W/m<sup>2</sup>) in the measurement of  $P''_{cool}(T)$  arising from the limitation of the tool and the power instability. We note that the fitted  $\varepsilon_{TARC,v}$  is the thermal emittance in the vacuum condition, and can be related to the ambient condition thermal emittance  $\varepsilon_{TARC,w}$  by

$$\varepsilon_{\text{TARC,v}} = \gamma \cdot \varepsilon_{\text{TARC,w}}$$

in which the coefficient  $\gamma$  is calculated to be about 0.7 from the spectra of TARC (Fig. S4), and is coincidentally the same for both the M state and the I state.

## Simulation of surface temperature

The stabilized temperature of a surface  $(T_s)$  with given solar absorptance (A) and thermal emittance  $(\varepsilon)$  was calculated based on adiabatic approximation, assuming negligible heat transfer between the

surface and the underlying structure. The key climate parameters for a specific city or region, including air temperature ( $T_a$ ), dew point temperature ( $T_d$ ), wind speed (v), solar irradiance (I) and cloud coverage factor (CF) are obtained from TMY3 weather files available from the U.S. Department of Energy (37). The thermal emittance of the TARC was set at 0.20 for  $T_s < 19$  °C and 0.90 for  $T_s > 27$  °C, and approximated by a linear interpolation in the transition region (19 °C  $\leq T_s \leq 27$  °C). Based on this setup, the all-year-around temperature map (Fig. 4C) of TARC and of all conventional materials with an arbitrary combination of static A and  $\varepsilon$  were calculated and compared. More details of the simulation can be found in the Supplementary Text, Note A, Section I.

## **Projection of energy savings**

In hour-of-year *i*, we define heating degrees  $D_{h,i} = (T_{set,heat} - T_{s,i})_+$  and cooling degrees  $D_{c,i} = (T_{s,i} - T_{set,cool})_+$ , where  $x_+ = x$  if x > 0, or 0 otherwise. The annually averaged heating degrees and cooling degrees are denoted by  $D_h$  and  $D_c$ , respectively.

Rosado & Levinson (38) simulated the annual space heating source energy savings  $\Delta S_h$  (typically negative) and the annual space cooling source energy savings  $\Delta S_c$  (typically positive) attained by increasing roof albedo for various categories and vintages of buildings in 15 U.S. climates zones and 16 California climate zones. All savings are normalized to roof area. Note that  $\Delta S_h$  and  $\Delta S_c$  are not directly presented by Rosado & Levinson but can be estimated from the heating, cooling, and fan energy uses reported in that work, as described in Appendix I of the Supplementary Text.

Summing  $\Delta S_h$  and  $\Delta S_c$  yields the annual space-conditioning (heating + cooling) source energy savings  $\Delta S$ . U.S. cool-roof space heating energy savings  $\Delta S_h$  and space cooling energy savings  $\Delta S_c$  reported by Rosado & Levinson are regressed against our own calculations of the reductions in annual average heating degrees  $\Delta D_h$  and cooling degrees  $\Delta D_c$ , respectively. These linear fits of the form  $\Delta S_h = \alpha_h \Delta D_h$  and  $\Delta S_c = \alpha_c \Delta D_c$  yield  $\Delta S = \alpha_h \Delta D_h + \alpha_c \Delta D_c$ . Specifically, energy simulations for four static roofing materials with  $\varepsilon_w = 0.90$  and solar reflectance R = 0.10, 0.25, 0.40, or 0.60 (A = 1 - R = 0.90, 0.75, 0.60, or 0.40) were selected for the extraction of the coefficients  $\alpha_h$  and  $\alpha_c$  using the material with A = 0.90 as the baseline.

To evaluate the potential space-conditioning source energy savings (SCSES) per unit roof area obtained by using TARC instead of a reference roofing surface with static solar absorptance  $A_{ref}$  and static thermal emittance  $\varepsilon_{ref}$ , we calculated, in each city,  $\Delta D_{h,TARC}(A_{ref}, \varepsilon_{ref}) \equiv D_{h,TARC} - D_h(A_{ref}, \varepsilon_{ref})$  and  $\Delta D_{c,TARC}(A_{ref}, \varepsilon_{ref}) \equiv D_{c,TARC} - D_c(A_{ref}, \varepsilon_{ref})$ , varying  $A_{ref}$  and  $\varepsilon_{ref}$  from 0 to 1. We then computed space-conditioning source energy savings  $\Delta S_{TARC}(A_{ref}, \varepsilon_{ref}) = \alpha_h \Delta D_{h,TARC}(A_{ref}, \varepsilon_{ref}) + \alpha_c \Delta D_{c,TARC}(A_{ref}, \varepsilon_{ref})$  for each building category and vintage of interest. The minimum  $\Delta S_{TARC}(A_{ref}, \varepsilon_{ref})$  (namely, SCSES<sub>min</sub>) for all possible existing roof coating properties is taken as the figure of merit for each combination of local climate, building category, and vintage group. In Fig. 1C and Fig. 4C, we used three different icons to represent three different types of competing materials where SCSES<sub>min</sub> occurs --- that is, the strongest rivals to TARC in energy saving. The triangle, circle, and square icons represent metallic, white, and dark roof coatings, with typical ( $A_{ref}, \varepsilon_{ref}$ ) parameters around (0.35, 0.25), (0.05, 0.95), and (0.95, 0.95), respectively.

The results in both Fig. 1C and Fig. 4C are based on the dominant resident building prototype in the U.S. (48), which is a single-family home built prior to 1980. Minimum annual source energy savings per unit roof area for single-family homes and apartment buildings built before 1980, between 1980 and 1999, and recently are presented in Table S4.

## **Supplementary Text**

## Note A. Simulations of surface temperature and energy saving advantage of TARC

## Section I. Calculation of surface temperature in adiabatic approximation

The stabilized surface temperature of TARC and any material with given solar absorptance (A) and sky-window thermal emittance ( $\varepsilon_w$ ) can be calculated based on the adiabatic approximation (49), assuming negligible heat exchange with the underlying structure. The thermal equilibrium is balanced by the three heat exchange components on the surface, which are net long-wave thermal radiative loss ( $q_{LW}$ ), short-wave solar absorption ( $q_{SW}$ ), and convective heat loss ( $q_c$ ), respectively. The equation for thermal equilibrium is denoted as

$$q_{\rm SW} - q_{\rm LW} - q_{\rm c} = 0 \tag{S1}$$

Each term can be expanded as

$$q_{\rm SW} = A \cdot I \tag{S2}$$

$$q_{\rm LW} = \varepsilon_{\rm w}(T_{\rm s}) \cdot \sigma \left(T_{\rm s}^4 - T_{\rm sky}^4\right) \tag{S3}$$

$$q_{\rm c} = h_{\rm c}(T_{\rm s} - T_{\rm a}) \tag{S4}$$

In these equations, *I* is the solar irradiance,  $T_s$  is the absolute surface temperature,  $T_{sky}$  is the absolute sky temperature,  $h_c$  is the convective heat transfer coefficient, and  $T_a$  is the absolute air temperature. The model schematic is shown in Fig. S16. We note that though Eq. (S3) is typically applied to broad-band IR emitters with uniform spectral emittance from 2 µm to 30 µm, it can be safely extended to the case of TARC. To support this argument, we calculated the net radiative cooling power for TARC and that of a broad-band emitter with constant emittance equal to  $\varepsilon_w$  of TARC. The general form of net radiative cooling power  $P_{net}$  is:

$$P_{\rm net} = P_{\rm rad} - P_{\rm atm} \tag{S5}$$

$$P_{\rm rad} = \iint B(\lambda, T_{\rm s})\varepsilon_{\rm s}(\lambda, \Omega, T_{\rm s})\cos\theta d\Omega d\lambda \tag{S6}$$

$$P_{\rm atm} = \iint B(\lambda, T_{\rm a}) \varepsilon_{\rm atm}(\lambda, \Omega) \varepsilon_{\rm s}(\lambda, \Omega, T_{\rm s}) \cos\theta d\Omega d\lambda \tag{S7}$$

in which  $P_{rad}$  and  $P_{atm}$  are radiative power of the surface and absorbed radiation power from the atmosphere, respectively; *B* is the black body spectral radiance;  $\varepsilon_s$  is the surface's spectral emittance; and  $\varepsilon_{atm}$  is the atmospheric spectral emittance. For simplicity, we assume that TARC is in M state when  $T_s \ge 25$  °C, and in I state when  $T_s \le 20$  °C. The  $\varepsilon_s$  of the M state and the I state can be found in Fig. S4, and  $\varepsilon_{atm}$  is retrieved from previous studies (50). The spectral integration range is 2-30 µm. The results are presented in Fig. S17, showing that the net radiative cooling power of TARC and the broad-band emitter are very similar for different scenarios, with deviation < 10% even for extreme conditions ( $T_s - T_a > 20$  °C). Therefore, the approximation of TARC as a broadband emitter in the calculation of long-wave radiative power ( $q_{LW}$ ) is valid.

The solar irradiance (*I*) and the air temperature ( $T_a$ ) can be directly imported from the TMY3 weather files included in the EnergyPlus climate database (*37*). For an infinitely large surface, the convective heat transfer coefficient ( $h_c$ , W/(m<sup>2</sup>·K)) is related to the wind speed (v, m/s) by (51)

$$h_{\rm c} = 5.6 + 5.1 \times \sqrt{\nu}$$
 (S8)

The sky temperature is approximated by

$$T_{\rm sky}^4 = \varepsilon_{\rm sky} \cdot T_{\rm a}^4 \tag{S9}$$

$$\varepsilon_{\rm sky} = \varepsilon_{\rm clear\_sky} + 0.8(1 - \varepsilon_{\rm clear\_sky}) \cdot \rm CF \tag{S10}$$

$$\varepsilon_{\text{clear sky}} = 0.754 + 0.0044 \times T_{\text{d}} \tag{S11}$$

following a clear-sky model (52) combined with cloud correction (53), and the sky emissivity is denoted as  $\varepsilon_{clear_sky}$  and  $\varepsilon_{sky}$  in each model, respectively.  $T_d$  is the dew point temperature (in °C) and CF is the cloud factor (ranging from 0 to 1; 0 for no clouds and 1 for full coverage), each of which is available in the climate database. The solar absorptance A and the thermal emittance  $\varepsilon$  are property parameters for the materials either pre-set (for reference materials) or obtained from experimental characterization (for TARC). Note that for TARC,  $\varepsilon_w(T_s)$  is set to 0.20 for  $T_s < 19$  °C and to 0.90 for  $T_s > 27$  °C, and approximated by a linear interpolation in the transition region (19 °C  $\leq T_s \leq 27$  °C). The only unknown parameter  $T_s$  can thus be calculated by solving equations (S1-S4).

Fig. S18 shows an example of the simulated surface temperature mapping over the year for Baltimore, MD. The map is divided by 24 hour-period in one day and 12 months over one year. Each block shows the hour-of-day average temperature in a given month. Each hourly surface temperature can be calculated from the weather data in the corresponding time slot and the material parameters. The generation of  $T_s$  mapping over one year can be applied beyond TARC to other conventional materials with fixed A and  $\varepsilon$ . This method can also be extended to any other regions provided that the local climate information is available.

### Section II. Energy saving advantage of TARC compared to conventional materials

The annual space heating source energy savings  $\Delta S_h$  and annual space cooling source energy savings  $\Delta S_c$  can be projected from the corresponding heating degrees reduction  $\Delta D_{h,TARC}(A_{ref}, \varepsilon_{ref})$  and cooling degrees reduction  $\Delta D_{c,TARC}(A_{ref}, \varepsilon_{ref})$ , as described in the main text and the Materials and Methods section. The space conditioning source energy saving (SCSES) can be obtained by summing  $\Delta S_{\rm h}$  and  $\Delta S_{\rm c}$ , and an example of SCSES mapped with  $A_{\rm ref}$  and  $\varepsilon_{\rm ref}$  (each ranging from 0 to 1) for Baltimore is plotted in Fig. S19. The dashed boxes in the right panel of Fig. S19 correspond to the parameters of currently available roof coatings. The minimum point within the dashed box regions denotes the material that leads to the minimum space conditioning source energy saving (SCSESmin) compared to all other roof coatings. In other words, for each city, the most energy-saving material (with fixed  $A_{ref}$  and  $\varepsilon_{ref}$ ) is first selected from all existing roof materials, and then compared with TARC. Thus, SCSES<sub>min</sub> is the additional source energy saving of TARC compared to that "existing best energy-saving material". SCSES<sub>min</sub> is taken as the figure of merit of TARC for the local climate and used to plot the mapping in Fig. 1C. In this example, the value for Baltimore is 22.4  $MJ/(m^2 \cdot y)$ , suggesting that would yield at least an annual source energy saving of 2.64 GJ over any other existing roof coating materials for a typical single-family building with a roof area of 118 m<sup>2</sup>. Note that both the example in Fig. S18 and the map in Fig. 1C are based on the building prototype of single-family house built prior to 1980, which is the dominant residential building type in the U.S. according to the 2015 Residential Energy Consumption Survey microdata (2). The extracted coefficients for this and for other building prototypes can be found in Tables S2 and S3. The calculated  $SCSES_{min}$  for all residential building prototypes (including those plotted in Fig. 1C) can be found in Table S4. The minimum coefficient of determination  $R^2$  when fitting  $\alpha_h$  and  $\alpha_c$  for all cases is about 0.98.

## Appendix I. Estimation of space-heating source energy saving $(\Delta S_h)$ and space cooling source energy saving $(\Delta S_c)$

Rosado & Levinson (38) report for residential and commercial buildings across California and the United States gas heating (*H*), air conditioning (*C*), and fan (*F*) annual source energy uses in the heating, ventilation, and air conditioning (HVAC) system. Some of the fan energy is used to deliver heated air, and some is used to deliver cooled air. To calculate space heating source energy use  $S_h$  and space cooling source energy use  $S_c$ , we split annual fan source energy use into annual heating fan source energy use  $F_h$  and annual cooling source fan energy use  $F_c$ :

$$F = F_{\rm h} + F_{\rm c} \tag{S12}$$

The simplest scheme is to assume that  $F_h$  is proportional to H and that  $F_c$  is proportional to C:

$$F_{\rm h} = \beta_{\rm h} \, H \tag{S13}$$

$$F_{\rm c} = \beta_{\rm c} C \tag{S14}$$

We use subscript 1 to refer to the baseline condition, such that:

$$F_{\rm h,1} = \beta_{\rm h} H_1 \tag{S15}$$

$$F_{\rm c,1} = \beta_{\rm c} C_1 \tag{S16}$$

$$F_1 = F_{\rm h,1} + F_{\rm c,1} \tag{S17}$$

Since  $\beta_h$  and  $\beta_c$  are assumed to be constant for each building prototype (combination of building category, vintage, and location), we have similar expressions for gas heating source energy savings  $\Delta H$ , air conditioning source energy savings  $\Delta C$ , and fan source energy savings  $\Delta F$  relative to the baseline:

$$\Delta F_{\rm h} = \beta_{\rm h} \,\Delta H \tag{S18}$$

$$\Delta F_{\rm c} = \beta_{\rm c} \, \Delta C \tag{S19}$$

$$\Delta F = \Delta F_{\rm h} + \Delta F_{\rm c} \tag{S20}$$

Combining Eqs. (S15) to (S20) yields:

$$\beta_{\rm h} = \frac{F_1 \,\Delta C - C_1 \Delta F}{H_1 \Delta C - C_1 \Delta H} \tag{S21}$$

$$\beta_{\rm c} = \frac{H_1 \Delta F - F_1 \,\Delta H}{H_1 \Delta C - C_1 \Delta H} \tag{S22}$$

The baseline refers to the condition where  $\varepsilon_w = 0.90$  and A = 0.90, and the source energy uses are  $H_1$ ,  $C_1$ , and  $F_1$ . The source energy savings  $\Delta H$ ,  $\Delta C$ , and  $\Delta F$  for conditions with other A values (0.75, 0.60, 0.40) are available in the savings database described in the Section 2.8 of Rosado & Levinson (38).

Finally, the space heating energy saving  $\Delta S_h$  and space cooling energy saving  $\Delta S_c$  can be calculated as:

$$\Delta S_{\rm h} = \Delta H + \Delta F_{\rm h} = (1 + \beta_{\rm h}) \,\Delta H \tag{S23}$$

$$\Delta S_{\rm c} = \Delta C + \Delta F_{\rm c} = (1 + \beta_{\rm c}) \,\Delta C \tag{S24}$$

## **Supplementary Figures**

1. Coating of the Si wafer by polyimide	2. Deposition of WVO <sub>2</sub> by PLD	3. Patterning of the WVO <sub>2</sub>	4. Evaporation of BaF <sub>2</sub> and Ag layers
polyimide Si wafer	WVO <sub>2</sub>		Ag BaF2
5. Attachment of the flexible substrate	6. Etching of the Si substrate	7. Removal of the polyimide layer	Finalized TARC
Scotch tape		Scotch tape Ag BaF <sub>2</sub> WVO <sub>2</sub>	0 1 cm 2 <u>З0 µт</u>

Figure S1. Schematics and pictures for the fabrication of TARC.

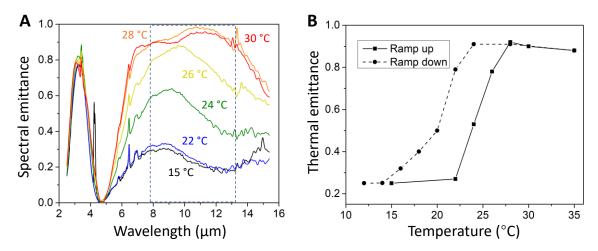
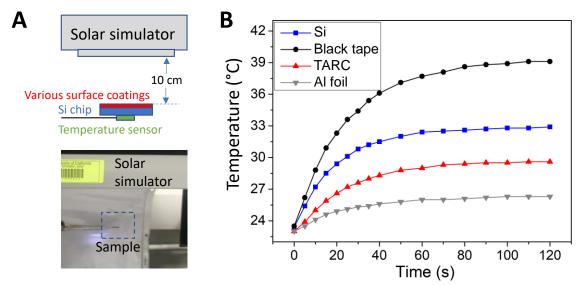
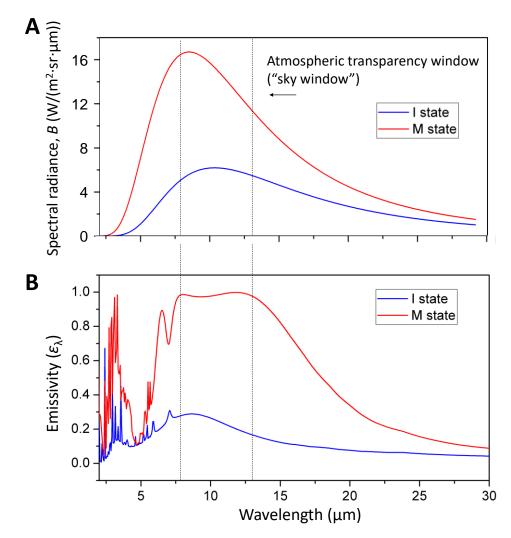


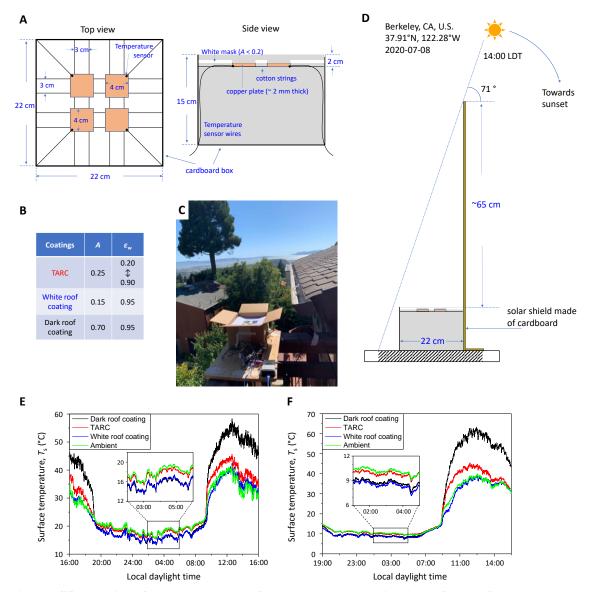
Figure S2. Detailed infrared spectra of TARC by Fourier transform infrared (FTIR) spectroscopy over the metal-insulator transition of  $W_x V_{1,x} O_2$  (*x*=1.5%). A. Spectral emittance of TARC at temperatures ranging from 15 °C to 30 °C (ramp-up branch). The dashed box bounds the sky window (8-13 µm). B. Sky-window emittance as a function of TARC surface temperature, including both the ramp-up and ramp-down branches.



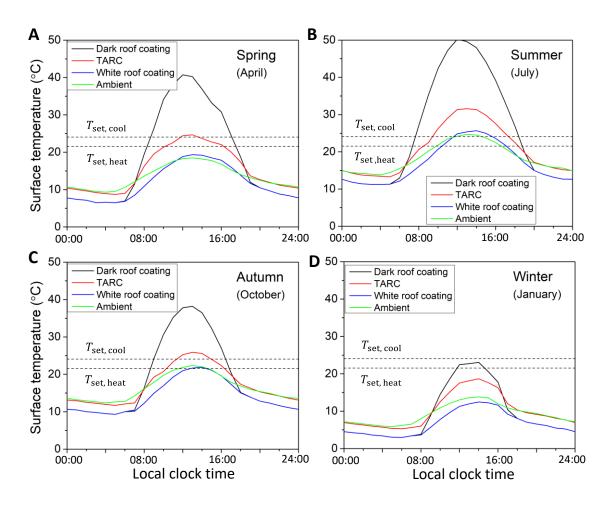
**Figure S3. Global solar absorptance characterization of TARC. A.** The schematic and a photo of the experiment setup. Various coatings were attached onto the top of a 1 cm  $\times$  1 cm Si chip, which was placed under a solar simulator (SS50B, Photo Emission Tech., Inc.). The temperatures were measured by a Pt temperature sensor. Before the measurement of each type of material surfaces, the shutter of the solar simulator remained closed, and the sample was first stabilized at the room temperature. Then the shutter was opened and the surface temperature was measured as a function of time. Only one sample was measured at one time, and this operation procedure was performed for four different surface coatings, maintaining the same setup geometry. The irradiance on the sample surface was 1 kW/m<sup>2</sup>. **B.** Measured temperature as a function of exposure time for different top coatings. The time when the shutter was opened is defined to be 0 s. The stabilized temperature is linearly related to the absorbed heat power and thus solar absorptance *A*. Based on the solar absorptance of other references ( $A_{Si} = 0.50$ ,  $A_{tape} = 0.95$ ,  $A_{AI} = 0.10$ ),  $A_{TARC}$  could be extracted as approximately 0.30, consistent with the theoretical prediction and other experimental results.



**Figure S4. Calibration of thermal emittance. A.** Spectral radiance of a blackbody at 0 °C and 60 °C. **B.** COMSOL-simulated spectral emittance of TARC in I state and in M state at wavelengths from 2 µm to 30 µm. The atmospheric transparency window is indicated by the grey dashed lines. Unlike the sky-window emittance, which is defined as  $\varepsilon_{\rm w} = (\int_{8\,\mu\rm m}^{13\,\mu\rm m} B(\lambda,T) \varepsilon(\lambda) d\lambda) / (\int_{8\,\mu\rm m}^{13\,\mu\rm m} B(\lambda,T) d\lambda)$ , the effective thermal emittance in vacuum environment is denoted as  $\varepsilon_{\rm vac} = (\int_{0}^{\infty} B(\lambda,T) \varepsilon(\lambda) d\lambda) / (\int_{0}^{\infty} B(\lambda,T) d\lambda) \approx (\int_{2\,\mu\rm m}^{30\,\mu\rm m} B(\lambda,T) \varepsilon(\lambda) d\lambda) / (\int_{0}^{\infty} B(\lambda,T) d\lambda)$  (28), where  $B(\lambda,T)$  is the spectral radiance of a blackbody. Based on the predicted spectral emittance of TARC, these two terms are related by  $\varepsilon_{\rm vac} = \gamma \cdot \varepsilon_{\rm w}$ , in which  $\gamma$  is about 0.7 for both the I state and the M state.



**Figure S5. Details of the outdoor performance characterization of TARC. A.** Schematic showing structure, geometry, and materials of the experimental setup. Identical 4 cm × 4 cm copper plates covered with different coatings were suspended by thin cotton strings (< 1 mm in diameter) in a cardboard box. A white paper mask was placed on top of the strings to minimize undesirable thermal artifacts from solar heating. B. Solar absorption and sky-window IR emittance of the measured material surfaces. **C.** A photo of the actual experiment setup. The measurement was carried out in an open-space rooftop balcony in Berkeley, CA (37.91°N, 122.28°W), with temperature time series recorded by a customized automatic temperature reading system. **D.** Schematic showing the solar shield used to block the direct solar radiation on the samples after 14:00 LDT. **E-F.** Additional experimental results measured on different days.



**Figure S6. Simulation of the surface temperature of a dark roof coating, TARC, and a white roof coating (the same three samples as those in Fig. 4) in four seasons in Berkeley, California (37.91°N, 122.28°W).** The ambient temperature data (green curves) are directly imported from the climate database (*37*) used for the simulation. The only condition where the white roof coating has an advantage over the TARC in energy saving is the daytime in summer (B), while for almost all other conditions (all night times, daytime in spring, winter, and most of autumn), the surface of temperature of TARC is closer to the ideal temperature zone between heating and cooling temperature setpoints (22 °C and 24 °C) than other roof coatings. Therefore, TARC saves more energy than the white roof coating from an all-season, day-and-night perspective.

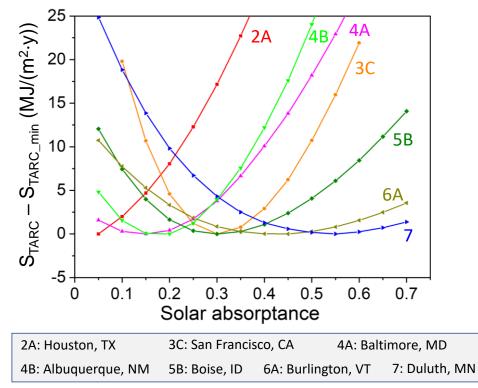


Figure S7. Calculation of space-conditioning source energy consumption of TARC ( $S_{TARC}$ ) in different cities as a function of the TARC solar absorptance (A). The indexes (2A, 3C, etc.) represent the different U.S. climate zones. In each curve,  $S_{TARC}$  is vertically offset by its minimum value for better display of the data in the plot. Each  $S_{TARC}$  curve is U shaped, suggesting that a solar absorptance neither too low nor too high is desirable. The ideal solar absorptance varies depending on the cities, and tends to shift to a higher value for those with colder climates, consistent with expectations. The experimental *A* of TARC is designed to be approximately 0.25 to 0.30, which is close to the optimal *A* (minimum point of the curves) for major U.S. cities.

Mean Space Conditioning Source Energy Saving (SCSES<sub>mean</sub>) of TARC Maximum Space Conditioning Source Energy Saving (SCSES<sub>max</sub>) of TARC

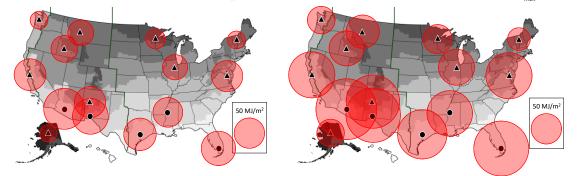
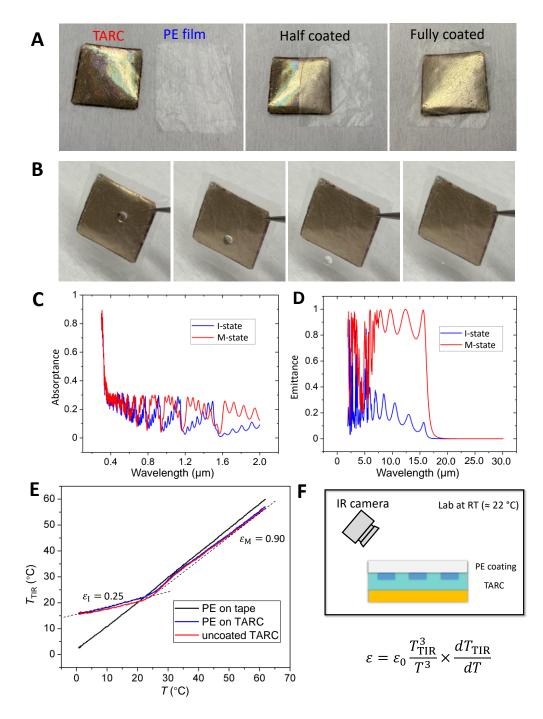


Figure S8. Mean space conditioning source energy saving (SCSES<sub>mean</sub>) and maximum space conditioning source energy saving (SCSES<sub>max</sub>) of TARC.  $SCSES_{mean}$  was calculated by averaging SCSEC over all available reference materials;  $SCSES_{max}$  is the maximum savings relative to all reference coatings.



**Figure S9. Demonstration of PE-coated TARC for practical applications. A.** Photos showing the setup of TARC coated with PE. **B.** Photos showing the hydrophobicity of PE-coated TARC, where water was dropped onto the PE-coated TARC sample. **C-D.** Calculated solar absorptance spectra and IR emittance spectra of TARC covered with PE. **E.** Experimental characterization of thermal emittance for a PE-coated TARC compared to an uncoated sample, showing that the additional PE coating has a negligible influence on the TARC thermal modulation performance. **F.** Schematic for the thermal emittance characterization via IR camera measurement.

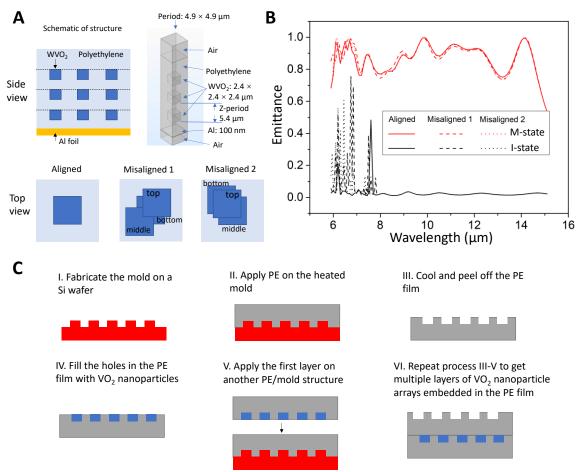


Figure S10. An alternative structural design and processing for mass production of TARC. A. Schematic design of the structure by tri-layer WVO<sub>2</sub> arrays embedded in polyethylene (PE) films. B. Spectral emittance of the I state and M state of the PE-based TARC film calculated by COMSOL. Significant emittance modulation ( $\Delta \varepsilon_w > 0.8$ ) is available, which also applies to structures with misaligned WVO<sub>2</sub> blocks in actual scenarios. C. Potential fabrication steps of the PE-based TARC film. Mass production is achievable via layer-by-layer nano-imprinting from a pre-fabricated Si wafer mold. Note that the Si wafer mold can be replaced by metallic imprinting mold for mass production in a roll-by-roll fashion.

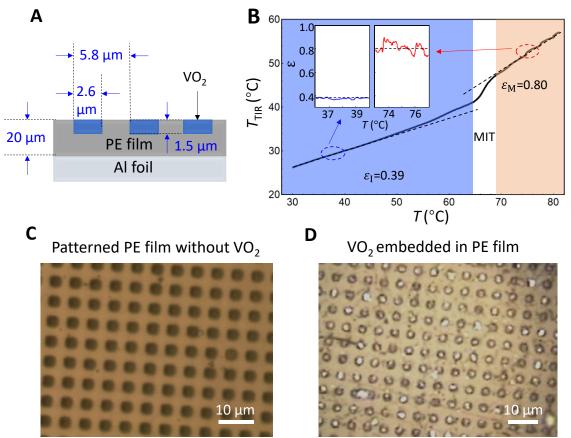


Figure S11. Preliminary results for the mass production of TARC based on the VO<sub>2</sub>embedded-in-PE approach, using only a single layer. The emittance (0.39) when VO<sub>2</sub> is in the I state can be further decreased by using thinner PE films. The emittance (0.80) when VO<sub>2</sub> is in the M state can be raised by stacking several PE-based TARC layers together. As a roof-coating material, PE-based coatings have a typical lifetime ranging from months (54) to years (55). This can be further prolonged by introducing advanced organic films (10,56). Furthermore, chemical additives and high-quality barrier coatings can also help optimize the durability of organic coatings (9).

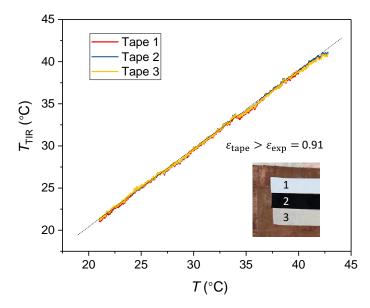
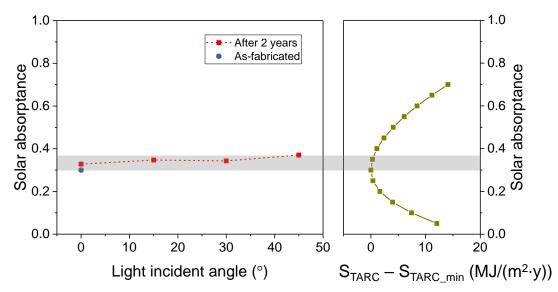
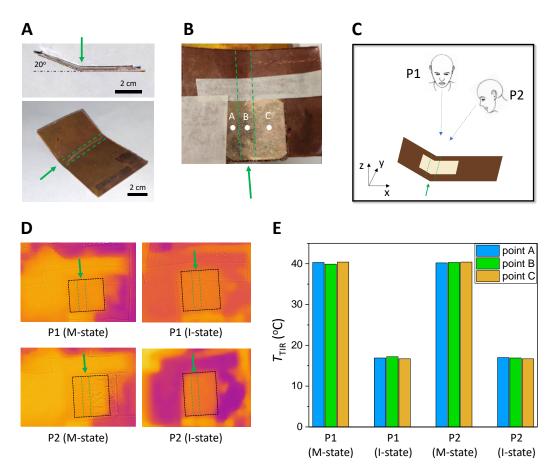


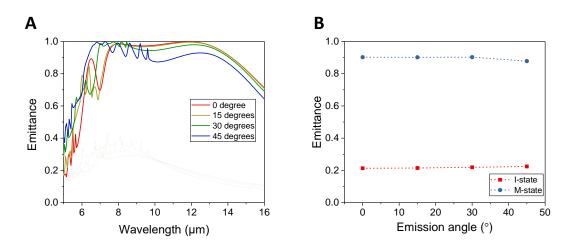
Figure S12. Experimental characterization of thermal emittance of the tape used in the experiments. Three pieces of tape were pasted onto a copper heating plate, and an IR camera (FLIR ONE) was used to measure the thermal IR temperature ( $T_{\text{TIR}}$ ) as a function of the contact temperature (T) measured by a Pt temperature sensor mounted near the tape. The thermal emittance of tape was extracted as 0.91. Due to heat convection and conduction, the tape's true temperature is always slightly lower than its measured contact temperature. Thus, the tape's true thermal emittance ( $\varepsilon_{\text{tape}}$ ) is higher than this extracted experimental thermal emittance ( $\varepsilon_{\text{exp}} = 0.91$ ), so the emissivity of tape is assumed to be 0.95 in other experiments according to reference (57).

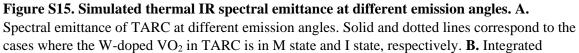


**Figure S13. Solar absorptance characterization of TARC at different light incident angles.** The experimental details and absorptance extraction method can be found in Fig. S3. The TARC sample has aged for two years before the angle testing. The solar absorptance at normal incidence increases only by 0.03 after the two-year aging. When the light incident angle increases from 0° (normal incidence) to 45°, the solar absorptance of TARC increases only by 0.04. This absorptance variation is compared to the normalized  $S_{TARC}$  in Boise, ID (details in Fig. S7), showing that the solar absorptance variation due to aging and light incident angles still falls well within the optimal range of  $S_{TARC}$ .



**Figure S14. TIR performance of TARC at different tilting angles and shape deformation. A.** Photos of a copper plate with a 20° bend. The bending edge is denoted by a green arrow, and the bending area is marked by two parallel green dashed lines (which are also shown in the other panels). **B.** A TARC was pasted onto the deformed copper plate, covering the bending edge. The three points (A, B, C) mark the positions where thermal IR temperatures ( $T_{TIR}$ ) are measured in panel E. Two pieces of tape applied next to TARC served as positioning markers in TIR images. **C.** Schematic for the experiment setup, where the default viewing angle of the IR camera (FLIR ONE) was set to 30°. The human-head icons show the IR camera positions and orientations in the YZ-plane (P1) and XZ-plane (P2). **D.** TIR images of the TARC (dashed black box) under different conditions. The copper plate was heated/cooled to about 40 °C/7 °C to test the device performance when WVO<sub>2</sub> is in M state/I state. Note that each TIR image has its individual color scale. **E.** Thermal IR temperature ( $T_{TIR}$ ) measured at different positions. The homogeneous  $T_{TIR}$  across the whole TARC indicate that its emittance is insensitive to both shape deformation and tilting angle.





emittance of TARC as a function of emission angles.

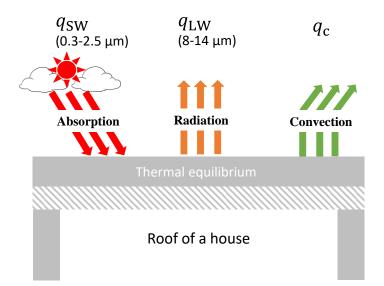


Figure S16. Schematic of the adiabatic approximation for surface temperature calculation. For a surface with given solar absorptance (*A*) and sky-window thermal emittance ( $\varepsilon_w$ ), this model enables the calculation of the surface temperature  $T_s$  in any climate.

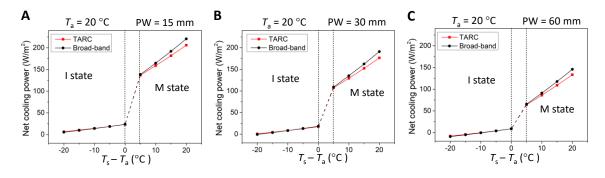


Figure S17. Calculated net radiative cooling power of TARC and a broad-band emitter with a constant emittance equal to  $\varepsilon_w$  of TARC. The results include conditions with various surface temperatures and precipitable water (PW). The atmospheric spectral emittances for different PW are obtained from (50).

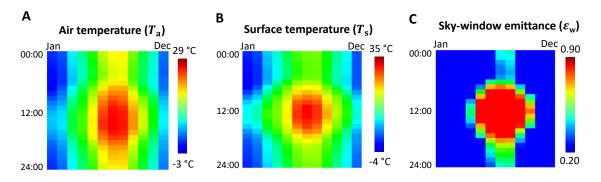
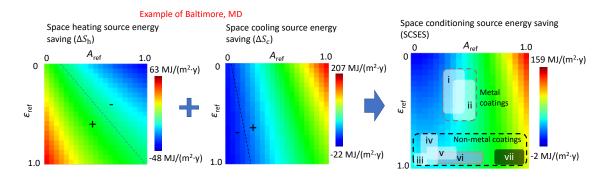


Figure S18. Example of surface temperature calculation based on the model for Baltimore, MD.



**Figure S19. Calculated energy saving advantage of TARC compared with conventional materials.** All savings are normalized to roof area. The indexes in the SCSES map represent: i, white metal plate; ii, aluminum pigments; iii, daytime radiative coolers; iv, silicone; v, acrylic; vi, polyurethane; and vii, asphalt.

## **Supplementary Tables**

Table S1. Detailed information for the compared thermal regulation technologies in Fig. 1B with references.

Category		Representative power (W)	Index in Fig. 1B	Reference
Peltier cooler	Peltier cooler		1	58-62
Refrigerator (pa	rts)	20-3,700	2	63-66
Caloric cooling		0.03-3,000	3	67-71
Membrane-assisted cooling		40-22,700	4	72-75
Air	Cooling	2,100-17,600	5	76.00
conditioner (with heating)	Heating	1,100-4,700	6	- 76-82
Joule heating		70-1,500	7	83-86
Gas heating		2,900-29,400	8	87-90

Part A. Energy-consuming devices<sup>a</sup>

## Part B. Energy-free materials

	Category	Representative power (W/m <sup>2</sup> )	Index in Fig. 1B	Reference
	Nanoparticles	70-93	9	9,91
	Wood	16	10	11
Daytime radiative	Polymer membranes	20-127	11	8,10,92-94
cooler <sup>b</sup>	Multilayer nanofilms	14-41	12	95,96
	Soft textiles (together with thermal emitters) <sup>c</sup>	60-96	13	39,97
TARC (this work)	Cooling (without sunlight)	0-110	14	This work
work)	Heating (with sunlight)	0-250	14	
Typical dark roof paint		600-980	15	98-100
Standard brick	Heating (with AM1.5	350-890	16	100-102
Al coating	sunlight)	120-300	17	100-102
Asphalt shingles		650-970	18	103-105

<sup>a</sup> Although most heaters/coolers can adjust their heating/cooling power from zero to the maximum power, in the master plot (Fig. 1B) we only graphed the ranges for the maximum heating/cooling powers.

<sup>b</sup> Only considers the net cooling power under direct sunlight in field experiments—i.e., cooling power minus heating power. Simulation data are not included here. If possible, the surface temperature is set to be ambient temperature.

<sup>c</sup> Soft textiles are usually tested while they are placed on thermal emitters. Details can be found in (39,97).

Climate zone	Location	Oldest (p	ore-1980)	Older (1980-1999)		New	
	Location	$lpha_{ m h}$	$\alpha_{\rm c}$	$\alpha_{\rm h}$	$\alpha_{ m c}$	$lpha_{ m h}$	$lpha_{ m c}$
1A	Miami, FL	31.48	24.52	21.24	10.55	18.87	8.53
2A	Houston, TX	33.00	22.61	23.30	11.28	19.93	7.67
2B	Phoenix, AZ	22.55	25.56	16.48	10.64	16.23	11.34
3A	Memphis, TN	17.97	23.29	11.77	11.37	11.71	11.25
3B	El Paso, TX	20.94	22.43	17.69	10.57	13.54	7.34
3C	San Francisco, CA	28.75	26.94	21.18	9.87	13.56	6.57
4A	Baltimore, MD	26.36	30.81	13.50	12.21	9.46	7.10
4B	Albuquerque, NM	20.90	27.17	14.64	10.45	13.26	7.72
4C	Seattle, WA	17.85	11.84	11.75	4.98	8.69	3.66
5A	Peoria, IL	18.19	22.76	11.13	10.44	8.85	7.47
5B	Boise, ID	21.86	16.86	12.56	6.55	11.85	4.81
6A	Burlington, VT	13.69	14.50	11.52	8.67	8.97	7.01
6B	Helena, MT	22.47	24.02	11.97	8.98	10.19	5.53
7	Duluth, MN	17.68	19.07	10.91	7.91	8.97	7.18
8	Fairbanks, AK	14.19	15.91	8.67	6.87	8.05	5.67

Table S2. Regressed coefficients  $\alpha_h$  and  $\alpha_c$  for single-family homes of different vintages in cities across the U.S. Units: MJ/(m<sup>2</sup>·y·K), where m<sup>2</sup> refers to the roof area.

Clim	Location	Oldest (p	ore-1980)	Older (1980-1999)		New	
ate zone	Location	$lpha_{ m h}$	$\alpha_{ m c}$	$lpha_{ m h}$	$lpha_{ m h}$	$lpha_{ m c}$	$lpha_{ m h}$
1A	Miami, FL	43.99	24.37	28.93	10.34	22.30	7.93
2A	Houston, TX	44.84	21.93	27.89	10.69	22.30	6.67
2B	Phoenix, AZ	39.24	24.59	24.22	9.96	24.26	9.90
3A	Memphis, TN	21.97	22.52	13.55	10.96	12.39	9.79
3B	El Paso, TX	34.66	21.03	23.23	9.95	17.57	6.12
3C	San Francisco, CA	40.35	23.43	24.37	8.45	14.29	5.09
4A	Baltimore, MD	29.82	29.51	14.40	11.72	9.21	5.60
4B	Albuquerque, NM	30.39	25.10	16.81	9.82	14.48	6.25
4C	Seattle, WA	21.17	10.77	12.22	4.48	7.86	3.65
5A	Peoria, IL	20.75	21.85	11.31	9.60	8.29	5.87
5B	Boise, ID	26.55	15.79	13.57	6.11	11.40	3.89
6A	Burlington, VT	14.48	12.37	12.68	9.76	7.87	5.05
6B	Helena, MT	25.71	22.05	12.37	7.67	8.97	3.86
7	Duluth, MN	21.36	20.78	11.44	8.83	8.20	5.61
8	Fairbanks, AK	17.45	15.28	9.55	6.43	7.42	4.53

Table S3. Regressed coefficients  $\alpha_h$  and  $\alpha_c$  for apartment buildings of different vintages in cities across the U.S. Units: MJ/(m<sup>2</sup>·y·K), where m<sup>2</sup> refers to the roof area.

		Singl	le-family ho	ome	Apartment building		
Climate zone	Location	Oldest (pre- 1980)	Older (1980- 1999)	New	Oldest (pre- 1980)	Older (1980- 1999)	New
1A	Miami, FL	-22.5	-6.6	-4.7	-17.5	-3.3	-2.4
2A	Houston, TX	12.4	9.5	8.6	25.1	11.9	7.18
2B	Phoenix, AZ	-2.1	9.7	8.4	24.1	22.3	22.4
3A	Memphis, TN	7.5	8.4	8.4	14.4	11.6	10.4
3B	El Paso, TX	20.2	16.4	11.3	32.7	15.0	9.0
3C	San Francisco, CA	14	4.2	2.8	10.6	2.03	1.8
4A	Baltimore, MD	22.4	8.5	4.9	21.1	8.1	3.7
4B	Albuquerque, NM	31.3	11.4	8.2	28	10.5	6.3
4C	Seattle, WA	3.1	0.83	-0.83	0.33	-2.1	-0.29
5A	Peoria, IL	14.5	5.7	4.6	13.9	5.9	3.45
5B	Boise, ID	15.1	5.6	3.9	13.7	5.0	2.14
6A	Burlington, VT	5.2	2.7	2.6	3.9	2.5	0.54
6B	Helena, MT	11.8	4.0	-0.039	10.0	1.6	-0.81
7	Duluth, MN	4.5	1.4	0.12	3.9	-0.031	-0.99
8	Fairbanks, AK	-1.0	-0.72	-2.7	-3.7	-3.4	-3.0

Table S4. Calculated  $SCSES_{min}$  as the first of merit for TARC in cities across the U.S. Units:  $MJ/(m^2 \cdot y)$ , where  $m^2$  refers to the roof area.

Climate	Location	Oldest (j	ore-1980)	Older (1980-1999)		New	
zone	Location	$eta_{ m h}$	$\beta_{\rm c}$	$\beta_{ m h}$	$\beta_{c}$	$eta_{ m h}$	β <sub>c</sub>
1A	Miami, FL	0.000	0.219	0.000	0.216	0.000	0.282
2A	Houston, TX	0.000	0.243	0.000	0.238	0.000	0.302
2B	Phoenix, AZ	0.078	0.218	0.114	0.207	0.106	0.209
3A	Memphis, TN	0.003	0.277	0.000	0.269	0.000	0.268
3B	El Paso, TX	0.068	0.262	0.074	0.253	0.091	0.299
3C	San Francisco, CA	0.122	0.395	0.124	0.411	0.151	0.395
4A	Baltimore, MD	0.038	0.344	0.026	0.341	0.045	0.385
4B	Albuquerque, NM	0.081	0.323	0.076	0.315	0.082	0.399
4C	Seattle, WA	0.045	0.486	0.040	0.493	0.054	0.507
5A	Peoria, IL	0.025	0.380	0.017	0.375	0.028	0.438
5B	Boise, ID	0.039	0.352	0.035	0.351	0.036	0.440
6A	Burlington, VT	0.030	0.465	0.070	0.213	0.034	0.517
6B	Helena, MT	0.042	0.434	0.040	0.425	0.039	0.538
7	Duluth, MN	0.086	0.211	0.085	0.222	0.057	0.499
8	Fairbanks, AK	0.060	0.215	0.060	0.213	0.060	0.212

Table S5. Regressed coefficients  $\beta_h$  and  $\beta_c$  (dimensionless) for single-family homes of different vintages in cities across the U.S.

Climate	Location	Oldest (pre-1980)		Older (1980-1999)		New	
zone	Location	$eta_{ m h}$	$\beta_{c}$	$\beta_{ m h}$	$\beta_{c}$	$eta_{ m h}$	β <sub>c</sub>
1A	Miami, FL	0.000	0.210	0.000	0.186	0.000	0.261
2A	Houston, TX	0.000	0.216	0.000	0.206	0.000	0.264
2B	Phoenix, AZ	0.121	0.207	0.207	0.196	0.205	0.196
3A	Memphis, TN	0.012	0.259	0.051	0.221	0.017	0.239
3B	El Paso, TX	0.075	0.249	0.088	0.237	0.119	0.283
3C	San Francisco, CA	0.115	0.372	0.113	0.349	0.154	0.335
4A	Baltimore, MD	0.033	0.333	0.035	0.297	0.118	0.295
4B	Albuquerque, NM	0.079	0.312	0.079	0.289	0.102	0.353
4C	Seattle, WA	0.043	0.462	0.040	0.416	0.075	0.384
5A	Peoria, IL	0.035	0.348	0.042	0.297	0.089	0.310
5B	Boise, ID	0.038	0.345	0.033	0.324	0.043	0.383
6A	Burlington, VT	0.060	0.345	0.058	0.310	0.072	0.362
6B	Helena, MT	0.044	0.422	0.065	0.305	0.069	0.370
7	Duluth, MN	0.073	0.373	0.073	0.326	0.075	0.419
8	Fairbanks, AK	0.057	0.309	0.061	0.225	0.061	0.224

Table S6. Regressed coefficients  $\beta_h$  and  $\beta_c$  (dimensionless) for apartment buildings of different vintages in cities across the U.S.

## **References and Notes**

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